

#### Generalized Slip Condition

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# IRREGULAR SURFACES, an equivalent description

#### Goal

A computationally cheap BC to simulate fluids flowing over generic rough walls



Rough surface under investigation, composed of a hexagonal periodic lattice:

- $\mathbb{FS}$  microscopic cell  $(l_1 \times l_2 \times l_3 \simeq l)$
- porosity  $\vartheta = |\mathbb{F}|/|\mathbb{FS}|$
- $\mathbb{ES}$  fictitious equivalent surface

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$$\epsilon = \frac{l}{L} \ll 1$$

# THE METHOD based on homogenization



Further hypotheses:

- up to  $\operatorname{Re} = U^{out}L/\nu = \mathcal{O}(1/\epsilon)$
- every kind of microscopic structure

After some algebra...



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#### THE RESULTING MODEL

$$u_{i} = -\epsilon^{2} \operatorname{Re} \underbrace{\mathcal{K}_{ij}}_{\partial x_{j}} \frac{\partial p}{\partial x_{j}} + \epsilon \underbrace{\mathcal{L}_{ilk}}_{\partial x_{k}} \left( \frac{\partial u_{l}}{\partial x_{k}} + \frac{\partial u_{k}}{\partial x_{l}} \right) |_{\mathbf{x} \in \mathbb{ES}},$$
  
$$\underbrace{\mathcal{K}_{ij}}_{\mathcal{K}_{ij}} = \langle \mathcal{K}_{ij} \rangle, \ \underbrace{\mathcal{L}_{ijk}}_{\mathcal{K}} = \langle \mathcal{L}_{ijk} \rangle, \ \mathcal{A}_{j} = \langle \mathcal{A}_{j} \rangle, \ \mathcal{B}_{lk} = \langle \mathcal{B}_{lk} \rangle, \ \langle f \rangle := \frac{1}{|\mathbb{FS}|} \int_{\mathbb{F}} f \ dV,$$

cf. Beavers & Joseph (1967) and Navier (1823).

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## THE RESULTING MODEL

$$u_{i} = -\epsilon^{2} \operatorname{Re} \frac{\mathcal{K}_{ij}}{\partial x_{j}} \frac{\partial p}{\partial x_{j}} + \epsilon \frac{\mathcal{L}_{ilk}}{\partial x_{k}} \left( \frac{\partial u_{l}}{\partial x_{k}} + \frac{\partial u_{k}}{\partial x_{l}} \right) |_{\mathbf{x} \in \mathbb{ES}},$$

$$\mathcal{L}_{ijk} = \langle \mathcal{L}_{ijk} \rangle, \ \mathcal{L}_{j} = \langle \mathcal{A}_{j} \rangle, \ \mathcal{B}_{lk} = \langle \mathcal{B}_{lk} \rangle, \ \langle f \rangle := \frac{1}{|\mathbb{FS}|} \int_{\mathbb{F}} f \ dV,$$

$$\begin{cases} \frac{\partial A_j}{\partial x_i} - \nabla^2 K_{ij} = \delta_{ij}, \\ \frac{\partial K_{ij}}{\partial x_i} = 0 \\ K_{ij} = 0 \quad \text{on } \partial \mathbb{S}, \\ \frac{\partial K_{jl}}{\partial x_3} + \frac{\partial K_{3l}}{\partial x_j} = 0, x_3 \in \mathbb{T} \end{cases}$$

$$\begin{cases} -\frac{\partial B_{j3}}{\partial x_i} + \nabla^2 L_{ij3} = 0, \\\\ \frac{\partial L_{ij3}}{\partial x_i} = 0, \\\\ L_{ij3} = 0 \quad \text{on } \partial \mathbb{S}, \\\\ \frac{\partial L_{pj3}}{\partial x_3} + \frac{\partial L_{3j3}}{\partial x_p} = \delta_{jp}, x_3 \in \mathbb{T}, \end{cases}$$

 $j = 1, 2, \forall i, K_{ii}, A_i, L_{ii3}$  and  $B_{i3} x_1, x_2$  periodic.

cf. Beavers & Joseph (1967) and Navier (1823).

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# MICROSCOPIC PROBLEMS: the surface permeability $\mathcal{K}_{ij}$

$$\begin{cases} \frac{\partial A_{j}}{\partial x_{i}} - \nabla^{2} \mathcal{K}_{ij} = \delta_{ij} \\ \frac{\partial \mathcal{K}_{ij}}{\partial x_{i}} = 0 \\ \mathcal{K}_{ij} = 0 \quad \text{on } \partial \mathbb{S} \\ \frac{\partial \mathcal{K}_{jl}}{\partial x_{3}} + \frac{\partial \mathcal{K}_{3l}}{\partial x_{j}} = 0, x_{3} \in \mathbb{T} \\ \mathcal{K}_{ij} \quad x_{1}, x_{2}\text{-periodic} \end{cases}$$

$$\mathcal{K}_{ij} = 0 \quad i \neq j \text{ by symmetry} \\ \mathcal{K}_{33} = 0 \text{ by axis orientation} \end{cases}$$
weak anisotropy:  $\mathcal{K}_{11} \neq \mathcal{K}_{22}$ 

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# MICROSCOPIC PROBLEMS: the surface permeability $\mathcal{K}_{ij}$



#### MICROSCOPIC PROBLEMS: the slip tensor $\mathcal{L}_{ijk}$

$$\begin{cases} -\frac{\partial B_{l3}}{\partial x_i} + \nabla^2 L_{il3} = 0\\ \frac{\partial L_{il3}}{\partial x_i} = 0\\ L_{il3} = 0 \quad \text{on } \partial \mathbb{S},\\ \frac{\partial L_{pl3}}{\partial x_3} + \frac{\partial L_{3l3}}{\partial x_p} = \delta_{lp}, x_3 \in \mathbb{T}\\ L_{il3} \quad x_1, x_2\text{-periodic} \end{cases}$$

 $\mathcal{L}_{ii3} 
eq 0 \quad i=1,2$  by symmetry and axis orientation



weak anisotropy:  $\mathcal{L}_{113} \neq \mathcal{L}_{223}$ 

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# MICROSCOPIC PROBLEMS: the slip tensor $\mathcal{L}_{ijk}$

$$\begin{cases} -\frac{\partial B_{l3}}{\partial x_i} + \nabla^2 L_{ll3} = 0 \\ \frac{\partial L_{ll3}}{\partial x_i} = 0 \\ L_{ll3} = 0 \quad \text{on } \partial \mathbb{S}, \\ \frac{\partial L_{pl3}}{\partial x_3} + \frac{\partial L_{3l3}}{\partial x_p} = \delta_{lp}, x_3 \in \mathbb{T} \\ L_{ll3} \quad x_1, x_2\text{-periodic} \end{cases}$$

$$\mathcal{L}_{il3} \neq 0 \quad i = 1, 2 \text{ by symmetry} \\ \text{and axis orientation} \end{cases}$$
weak anisotropy:  $\mathcal{L}_{113} \neq \mathcal{L}_{223}$ 

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#### VALIDATION: flows past rough spherical particles

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Free uniform unidirectional flow (U^{out}, 0, 0) past a spherical particle such that:
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- $\vartheta = 0.60$
- 1440 protrusions, i.e.  $\epsilon \approx 0.03$
- Re = 100
- $200R \times 80R \times 80R$
- 15 to 25 millions cells
- OpenFOAM in parallel with up to 500 cores
- $\bullet~$  up to  $1.2\times10^5~CPU$  hours for each DNS to reach the steady state at Re= 100



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# VALIDATION: local comparisons



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# VALIDATION: local comparisons



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# CONCLUSIONS and more

A robust BC for incompressible fluid flows above rough (and smooth?) surfaces has been developed and validated in a non trivial case (showing a good agreement also when the homogeneity is weakened)

- go further in the  $\epsilon$ -expansion to develop a strategy to better locate the equivalent surface in the macroscopic analogy
- extend the procedure in a stochastic direction to analyze irregular surfaces without a fixed shape of the protrusions
- study the dynamics of Janus-like spheres optimizing distribution and shape of the protrusions to enhance their dynamic performances
- following the path of drag reduction, homogenization of two phase flows is necessary

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